

RATES FOR EXCITATION OF THE CO₂ FERMI RESONANCE MEMBERS IN RF ELECTRIC FIELD

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Abstract. Rate coefficients for electron impact excitation of the members of the Fermi resonance of CO₂ in non-equilibrium conditions were calculated. Monte Carlo simulation of electron transport through gas in the presence of radio frequency electric field was employed in order to determine proper electron energy distribution functions (EEDF) for the given parameters – 100 MHz field frequency and the reduced field magnitude of 100 Td.

1. INTRODUCTION

CO₂ gas is extensively used in various devices that generate plasma by radio frequency discharge (Spencer et al. 2011, Srivastava et al. 2010). Electron collisions with CO₂ gas molecules occur frequently inside those devices and influence the further molecular dynamics. Vibrational excitation has great capacity to store energy which is later transferred between excited molecules. This is being used to achieve energy efficient dissociation, which is a key process in CO₂ conversion to ecologically friendly fuels. Rate coefficients are important input data in modeling conditions in discharge devices.

Fermi resonance was first recognized by Enrico Fermi in 1931 (Fermi 1931) as quasi degeneration of two vibrational levels, by which he explained one line in Raman spectra of CO₂, instead of two as then existing theoretical models predicted. Fermi resonances play important role in vibrational energy transfer. CO₂ molecule possesses two groups of nearly degenerate vibrational levels. The first one, called Fermi dyad, includes the lowest symmetrical stretch mode, (100), and the bending mode (020). The second one, named Fermi triad, consists of three modes – (200), (120) and (040). In the present work we are going to show time resolved rate coefficients for electron impact excitation of all these Fermi resonance members, which were obtained by means of a Monte Carlo simulation that treats electrons moving in gas filled space, under the action of a spatially uniform radio frequency

electric field. We chose the reduced electric field, E_R/N (root mean square value), to have magnitude of 100 Td. At this E_R/N value period averaged mean electron energy reaches 3.55 eV.

2. MONTE CARLO SIMULATION

The simulation follows up to 10^7 electrons, which are accelerated by the action of the radio-frequency electric field and subsequently decelerating upon collisions with CO_2 gas molecules, distributed in an infinite space. Electrons are treated one by one and their motion is determined in small time steps by numerical solution of the differential equations of motion. As soon as the steady state is reached, electron energy distribution function is repeatedly sampled, along with other relevant transport parameters, inside one cycle of the electric field and then averaged to acquire better statistics.

Data read from the input cross section database are used to calculate collision probabilities. A collision event is simulated by calling a random number generator.

A significant part of the cross sections for vibrational excitation in our database makes data obtained by coauthors of this paper in the past (Poparić et al. 2010). Those are cross sections for excitation of symmetrical stretch modes (100), (200) and so on up to (800). As it was said in the introduction, Fermi resonance consists of nearly degenerate vibrational modes. Separate contributions of the modes (100) and (020) were obtained after comparing energy loss spectra measured by (Poparić et al. 2010) with spectra reported by (Johnstone et al. 1995) and (Kitajima et al. 2001), whereby their peak deconvolution method and their angular distributions were applied. Cross sections for excitation of (120) and (040) are calculated values originating from (McCurdy et al. 2003).

Since simulation performed for relatively low reduced electric field, superelastic deexcitation from (010) * level to the ground (000) state, which is occurring with non-negligible probability at room temperature (Kato et al. 2008), was taken into account. Cross sections for this process were calculated by using the Klein-Rosseland formula.

2. RESULTS

The simulation shortly described in the previous section was prepared for the electric field parameters of 100 Td ($1 \text{ Td} = 10^{-21} \text{ Vm}^2$) and 100 MHz. The gas pressure was set to 1 Torr (133.3 Pa). EEDFS obtained within one period of oscillation by the simulation were used to determine rate coefficients. Rate coefficients were calculated by applying the formula:

$$K(\langle \varepsilon \rangle_t, t) = \sqrt{2/m} \int_{\varepsilon_{th}}^{\infty} \sigma(\varepsilon) \sqrt{\varepsilon} f_e(\langle \varepsilon \rangle_t, \varepsilon, t) d\varepsilon \quad (1).$$

In the equation above $\langle \varepsilon \rangle_t$ and $f_e(\langle \varepsilon \rangle_t, \varepsilon, t)$ represent the mean electron energy and the normalized EEDF, respectively, at the moment of time t , ε is the electron energy, $\sigma(\varepsilon)$ is the cross section for excitation of the given vibrational mode, while the threshold energy is denoted by ε_{th} .

Time resolved rate coefficients for excitation of (100) and (020) modes, belonging to Fermi dyad of CO₂, are shown in figure 1. The results obtained for excitation of the members of Fermi triad, namely of the (200), (120) and (040) modes, are shown in figure 2.

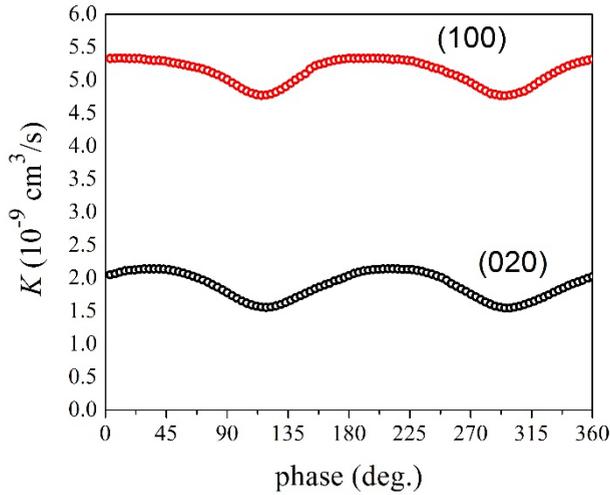


Figure 1: Rate coefficients for excitation of the Fermi dyad members.

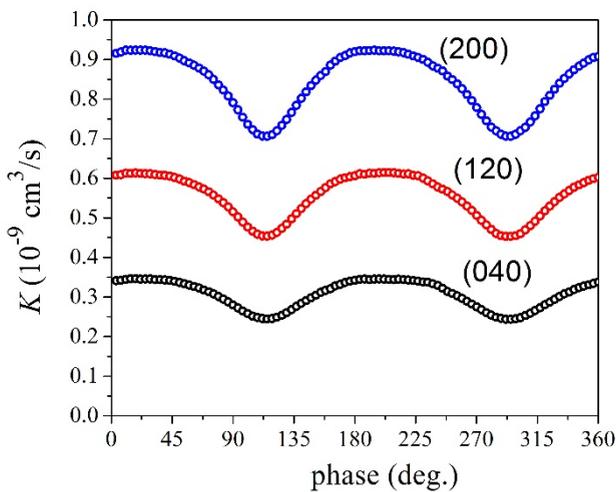


Figure 2: Rate coefficients for excitation of the Fermi triad members.

The obtained results indicate that rates for (100) and (020) are slightly shifted in phase, whereas all shown excitation rates for the members of Fermi triad are in phase. Another difference one may observe when comparing these two figures is that the amplitude of the members of Fermi dyad, shown in figure 1, is the same, but in case of Fermi triad (figure 2) amplitudes differ in magnitude for all three members. The intensity arrangement for these modes follows from the magnitudes of the corresponding cross sections. The time modulation of all these rates can be understood after performing detailed analysis of the position of the effective cross section function for the given mode with respect to the EEDF at the specific moment in time.

Acknowledgments

This work is partly supported by the Ministry of Education, Science and Technological Development of the Republic of Serbia.

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